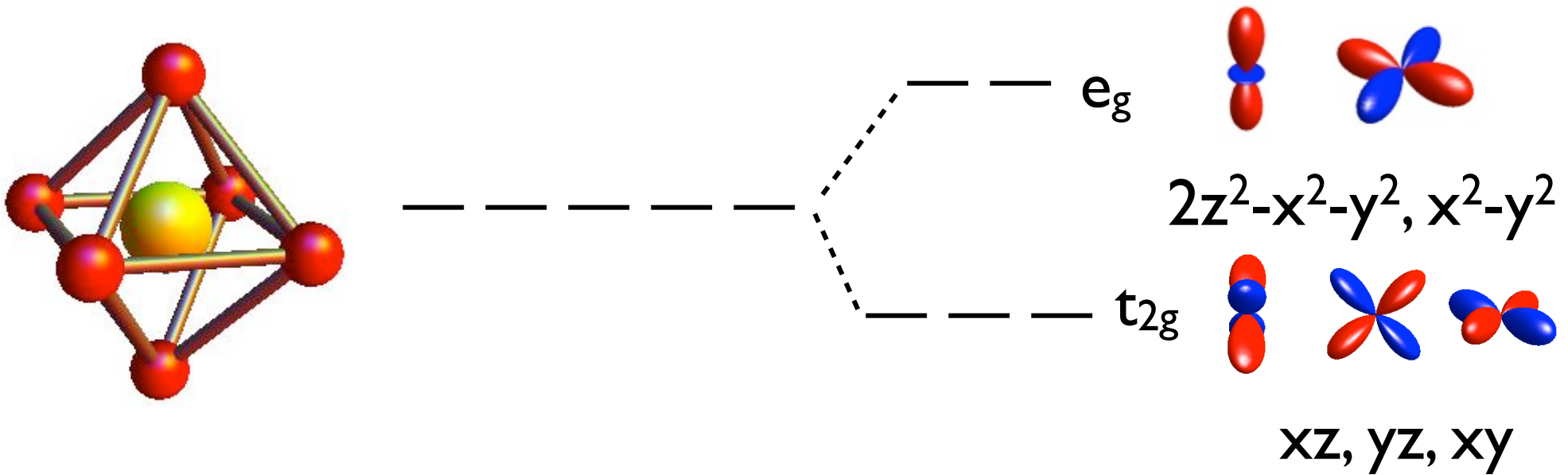


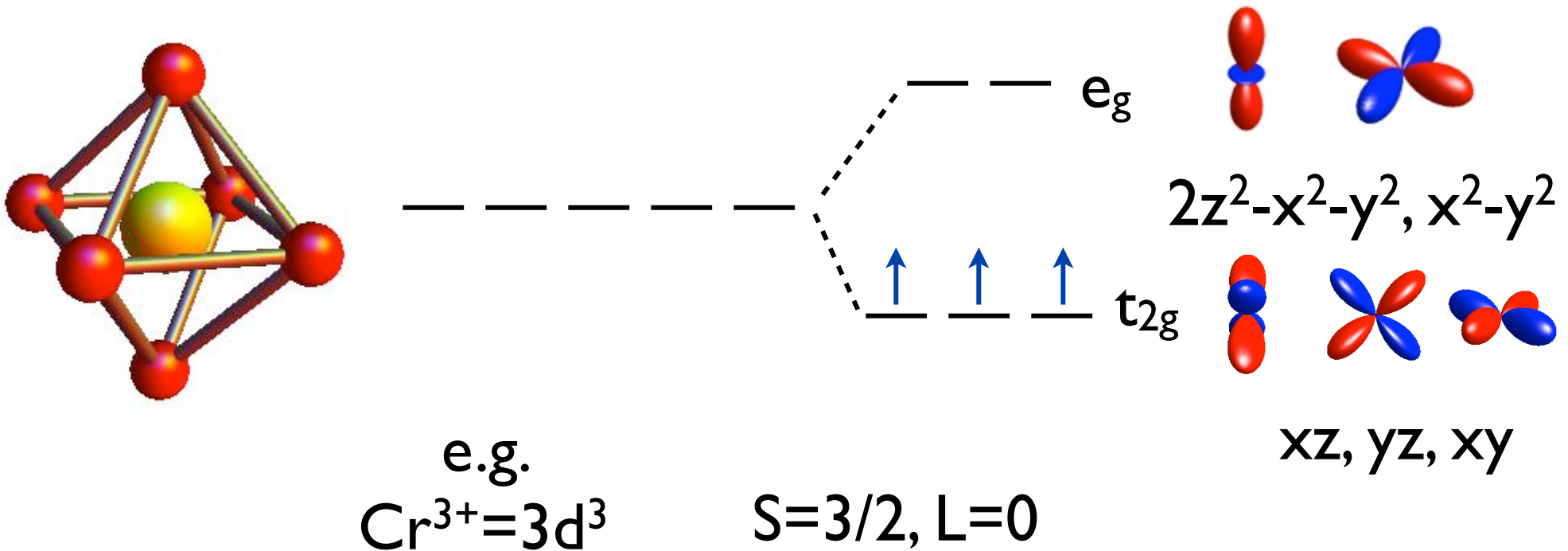
# Moments in Solids

- Example: cubic crystal field for a metal atom in an oxygen octahedron



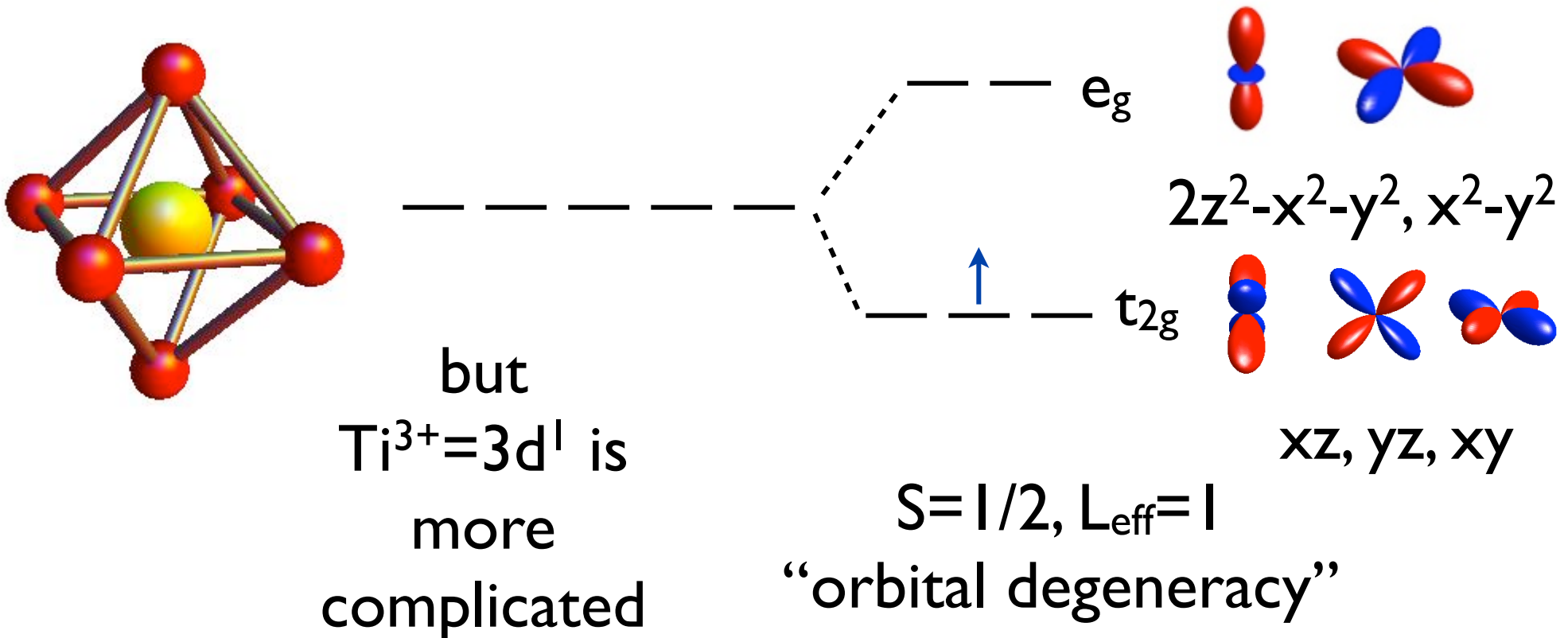
# Moments in Solids

- Example: cubic crystal field for a metal atom in an oxygen octahedron



# Moments in Solids

- Example: cubic crystal field for a metal atom in an oxygen octahedron



# Local moments

- Local moments are *not* part of band theory
  - Works in materials where electrons are localized to atoms, and delocalization is prevented somehow -- insulators
  - Such materials, which have partially filled shells but are insulating, are called Mott insulators
- How do we know they exist?

# Curie Susceptibility

- Existence of local moments means degenerate states
- By application of a small magnetic field, this degeneracy is split and a particular spin state is selected
- Expect large susceptibility  $\chi = \partial M / \partial H|_{H=0}$

# Curie Susceptibility

- Magnetic moment in general is proportional to spin

$$\mu = g\mu_B \mathbf{S} / \hbar$$

spin  $S$  quantum spin  
 $S^2 = S(S+1)\hbar^2$

g-factor  
 (could be a tensor)  
 $g \approx 2$  for pure spin moment

Bohr magneton  

$$\mu_B = \frac{e\hbar}{2m_e} = 9.3 \times 10^{-24} \text{ J/T}$$

$$= 0.671 \text{ K/T}$$

- Magnetic dipole interaction

$$\mathcal{H} = -\boldsymbol{\mu} \cdot \mathbf{H}$$

# Curie Susceptibility

- Thermodynamics  $M = -\frac{\partial F}{\partial H}$   $F = -kT \ln Z$

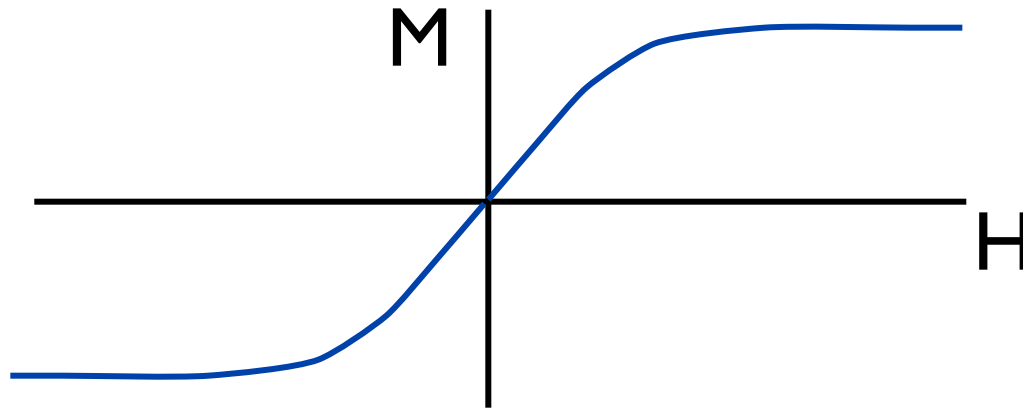
$$Z = \sum_n e^{-E_n/kT}$$

- For example,  $S=1/2$   $E_{\uparrow/\downarrow} = \mp g\mu_B H/2$

$$Z = 2 \cosh \left( \frac{g\mu_B H}{2kT} \right) \quad M = \frac{g\mu_B N}{2} \tanh \left( \frac{g\mu_B H}{2kT} \right)$$

# Curie Susceptibility

- Magnetization  $M = \frac{g\mu_B N}{2} \tanh\left(\frac{g\mu_B H}{2kT}\right)$



- Susceptibility

$$\chi = \frac{(g\mu_B)^2 N}{4k_B T} = \frac{A}{T}$$

Curie Law:  
indicates local moment

# Curie Susceptibility

- Thermodynamics  $M = -\frac{\partial F}{\partial H}$   $F = -kT \ln Z$

$$Z = \sum_n e^{-E_n/kT}$$

- General result

$$\chi = \frac{N(g\mu_B)^2}{3} \frac{S(S+1)}{kT}$$

plot  $\chi^{-1}$  versus T to extract “effective moment”

# Contrast with metals

Curie Law

$$\chi = \frac{N(g\mu_B)^2}{3} \frac{S(S+1)}{kT}$$

Pauli paramagnetism

$$\chi = V \frac{(g\mu_B)^2}{4} D(\epsilon_F)$$



basically  $kT \rightarrow \epsilon_F$

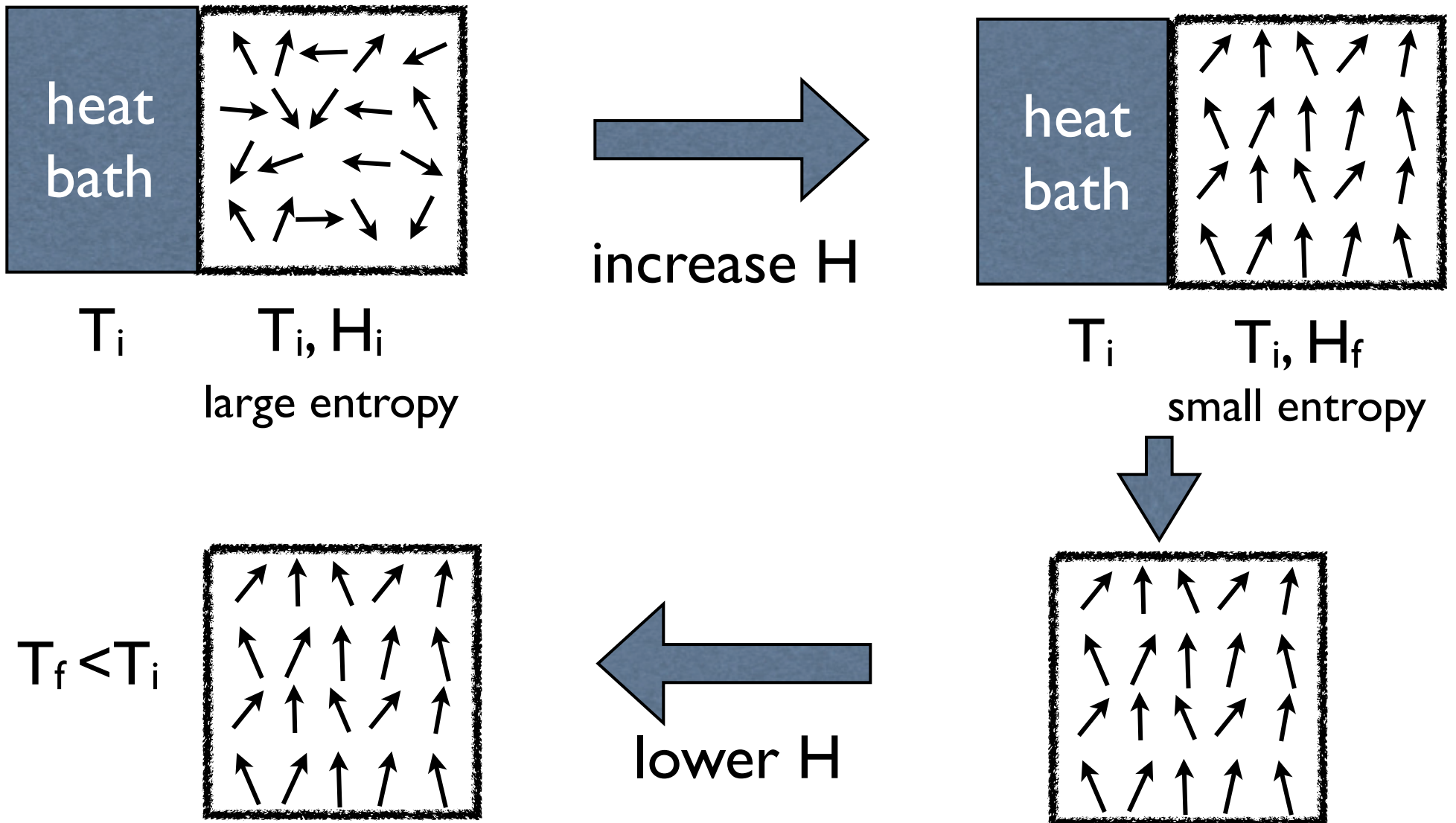
Does  $\chi$  really diverge for local moments?

**NO**, because they interact

# Magnetic cooling

- The large susceptibility of free spins at low temperature means they are easily aligned by small magnetic fields
- This alignment corresponds to a drastic reduction of entropy. One can use this control over entropy to remove entropy from another system, thereby cooling it.

# Magnetic Cooling



# Magnetic Cooling

- A → B: isothermal step - raise field, lower entropy
- B → C: adiabatic step - lower field, same entropy: lower temperature
- For paramagnetic spins,  $S = S(H/T)$ 
  - Hence  $H_1/T_f = H_2/T_i$

