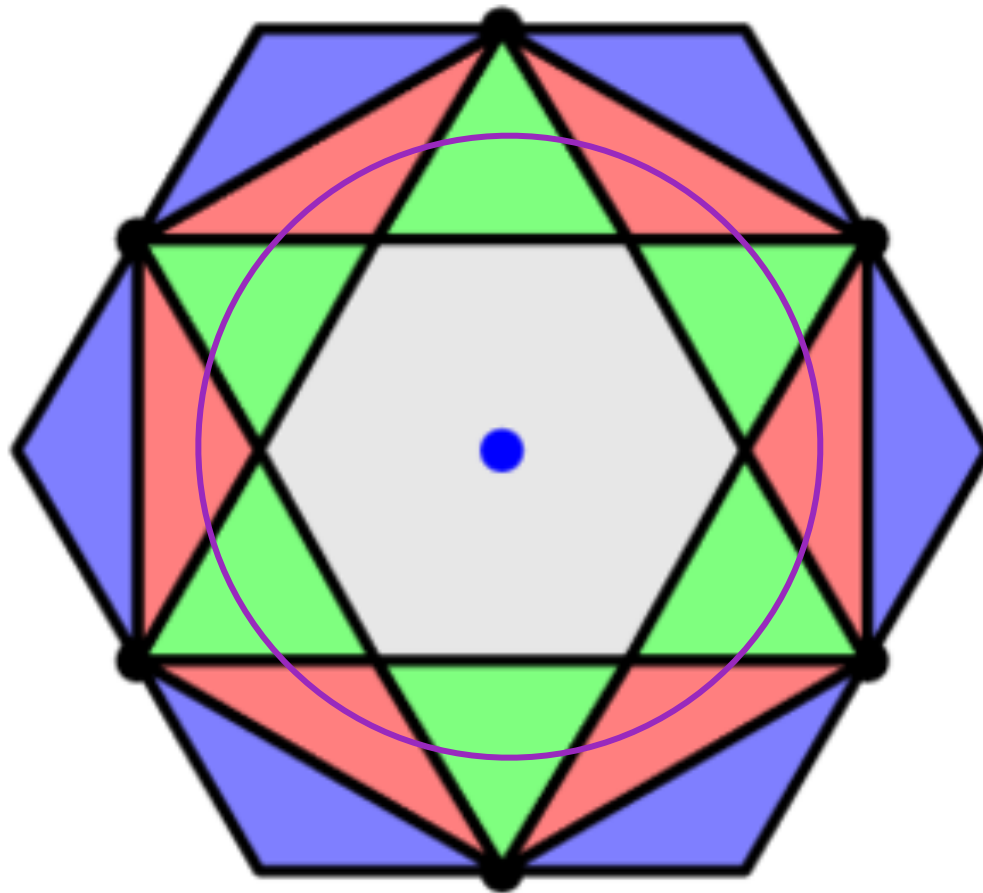
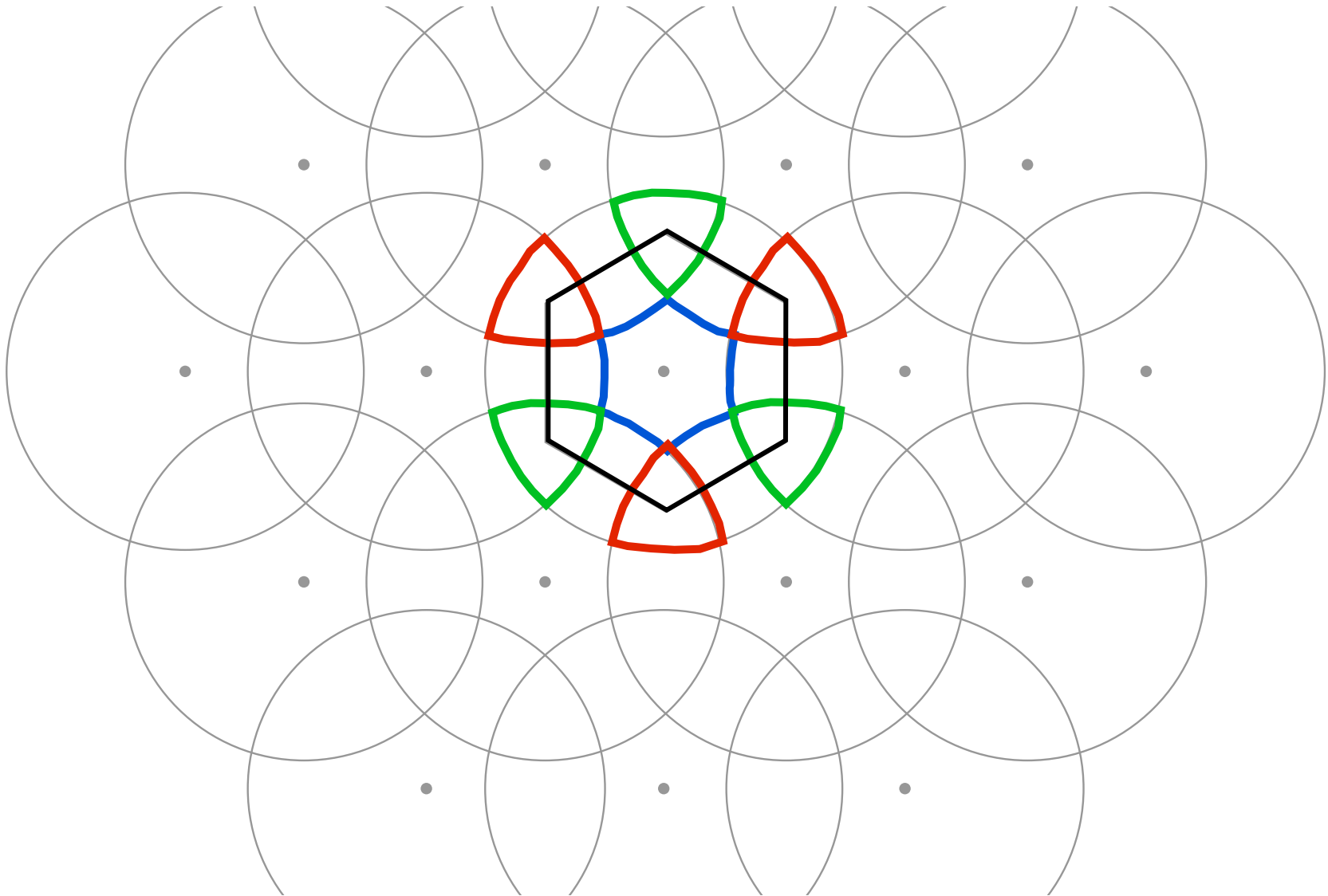


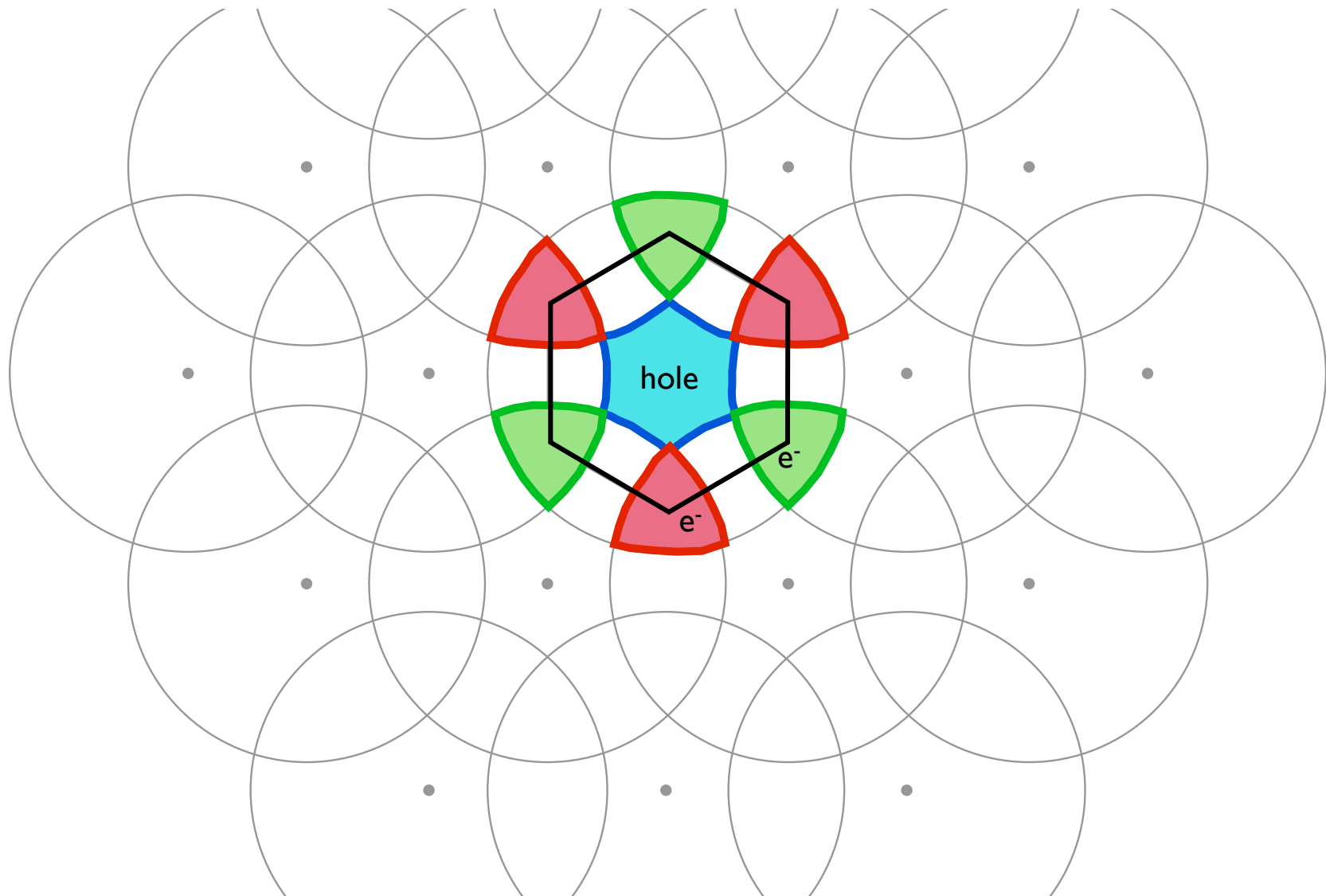
# Hexagonal BZs



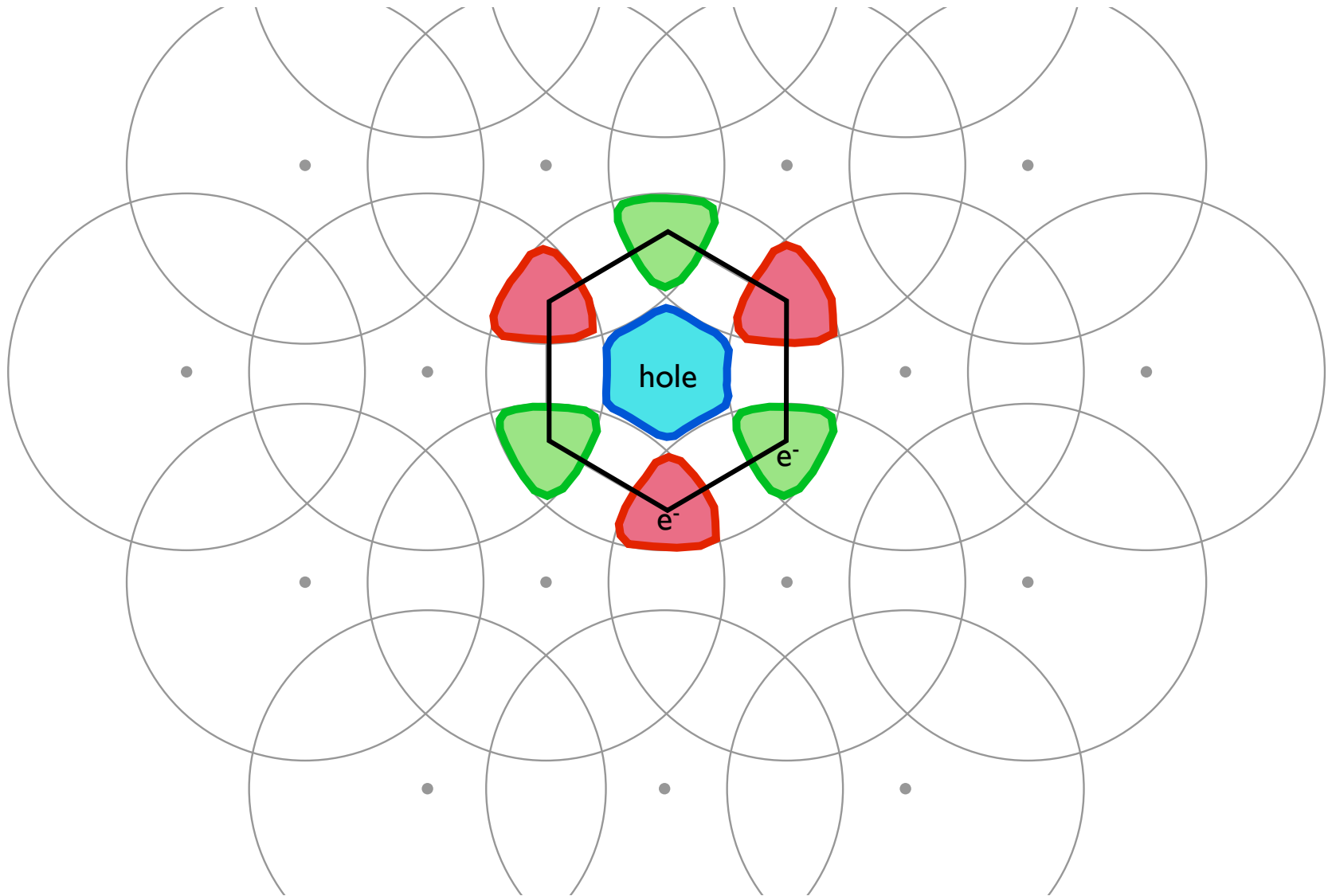
# Fermi surface pockets



# Fermi surface pockets



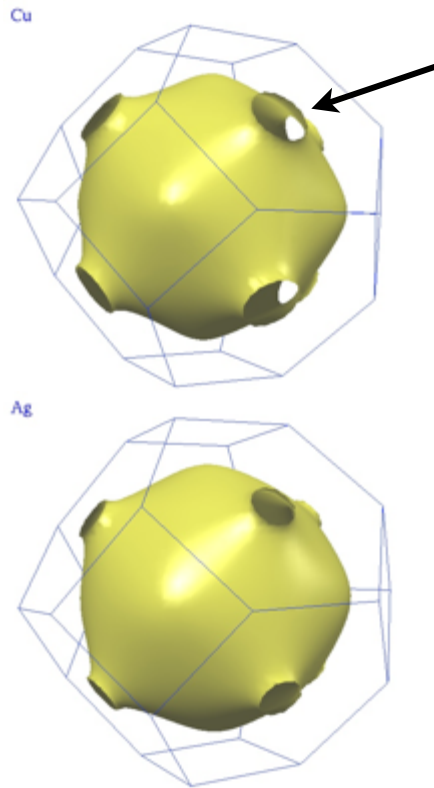
# w/ periodic potential



# Lessons

- If Fermi sphere extends past first Brillouin zone, several pieces of Fermi surface will exist in multiple bands in free electron limit
- Once weak interactions are added, “cuspy” Fermi surfaces become rounded, and usually do not intersect

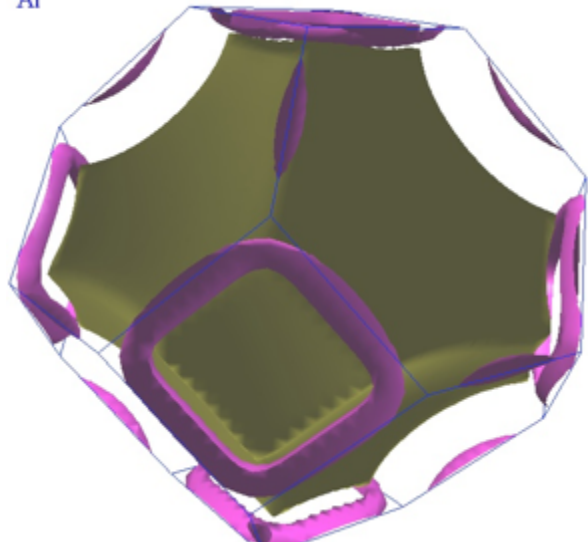
# Free electron-like FSs



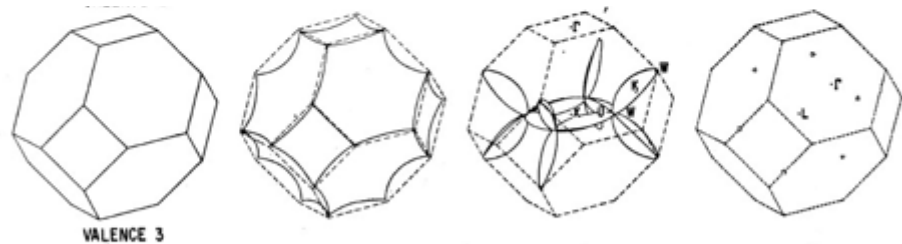
point inside the “neck” is on Bragg plane.  
Band gap opened at this  $k$ , lowering the energy of valence band so it moved inside the Fermi surface. This explains tendency to “necking”

# Free electron-like FSs

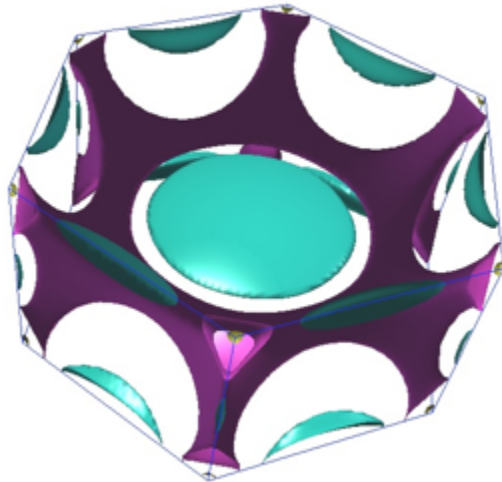
Al



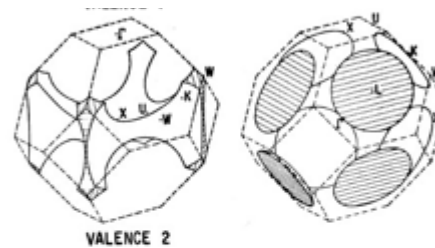
apparently, at least the larger Fermi surface of Al is free-electron-like.



Mg



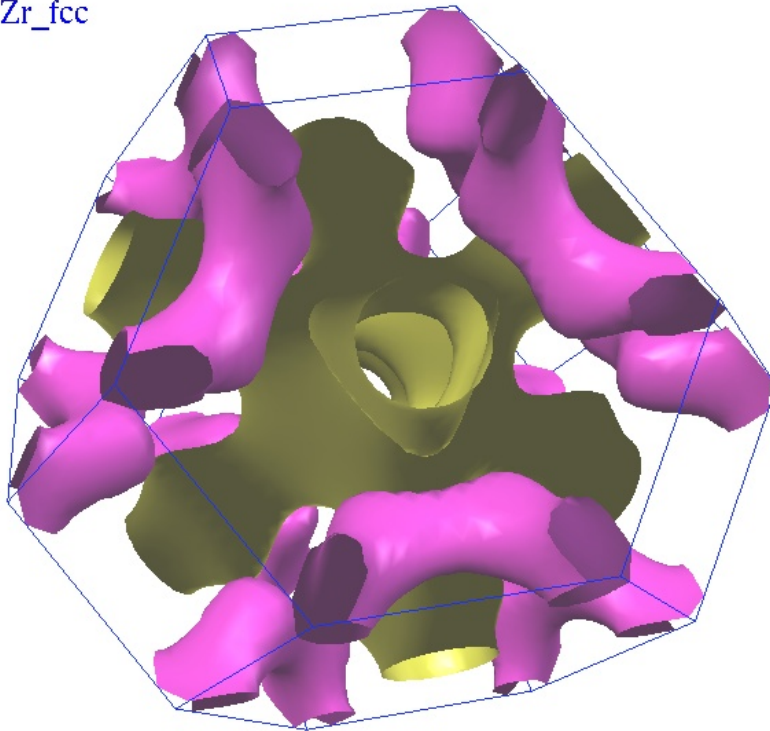
compare to Ch. 9, Fig. 1



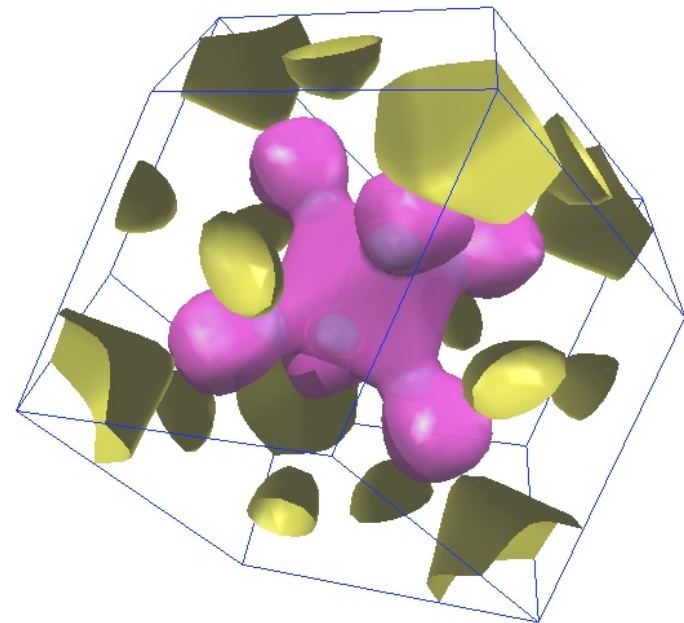


# Wild stuff

Zr\_fcc



Cr



# Tight binding method

- The opposite limit from NFEA - assume the ionic potential strongly confines electrons
  - only a small number of atomic orbitals are important
- We can try to construct Bloch states from these orbitals only
  - conceptually similar to making “bonding” and “anti-bonding” orbitals on molecules
  - but with  $10^{23}$  atoms instead of 2!

# Tight Binding

- Write the wavefunction as a superposition

$$\psi(r) = \sum_R \psi_R \phi(r - R)$$

↑ amplitude                      ← orbital of atom at R

- Amplitudes obey “discrete Schrodinger equation”

$$\hat{H}\psi_R = \epsilon_0 \psi_R - \sum_{R'} \gamma_{R,R'} \psi_{R'} = \epsilon \psi_R$$

atomic energy                      “hopping” amplitude -  
decays rapidly with distance

# Tight binding

- Often we assume just nearest-neighbor hopping
- Example: one dimensional chain

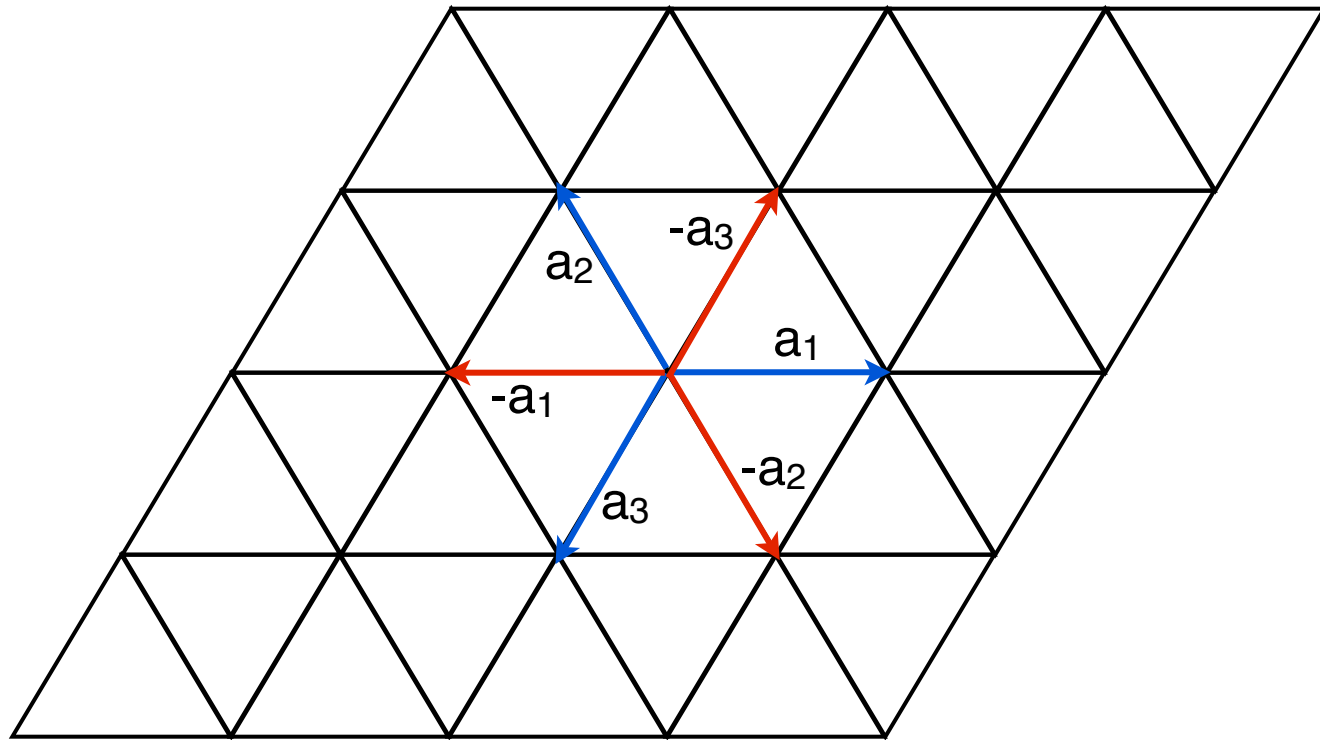
$$\hat{H}\psi_x = \epsilon_0\psi_x - \gamma(\psi_{x+a} + \psi_{x-a}) = \epsilon\psi_x$$

- Solve it?  $\psi_x = \bar{\psi}e^{ikx}$

$$\epsilon(k) = \epsilon_0 - 2\gamma \cos ka$$

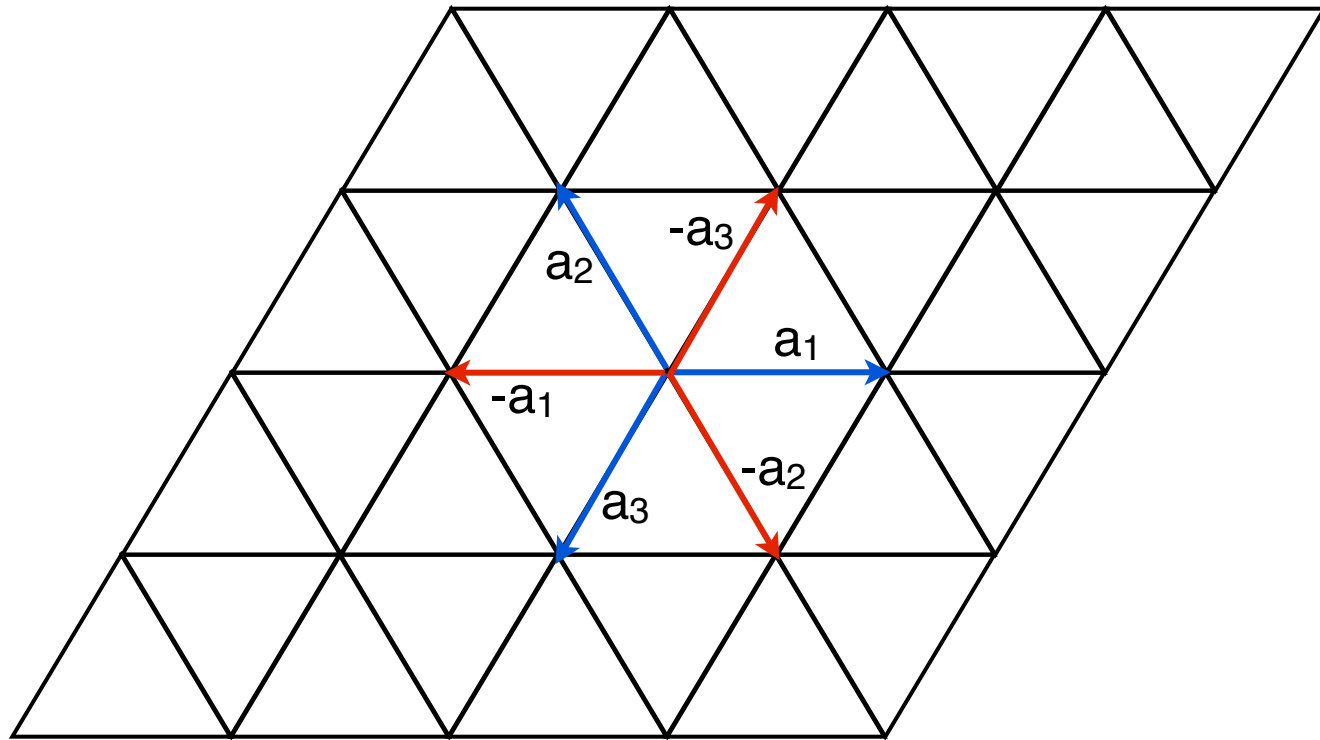
- Easily generalized to 2d and 3d lattices (see Kittel)

# 2d hexagonal lattice



$$\epsilon_0 \psi_R - \gamma \sum_{\langle R', R \rangle} \psi_{R'} = \epsilon \psi_R$$

# 2d hexagonal lattice



$$\epsilon_0 \psi_R - \gamma \sum_{i=1}^3 (\psi_{R+a_i} + \psi_{R-a_i}) = \epsilon \psi_R$$

# 2d hexagonal lattice

- Spectrum:

$$\epsilon(k) = \epsilon_0 - 2\gamma \sum_{i=1}^3 \cos \mathbf{k} \cdot \mathbf{a}_i$$

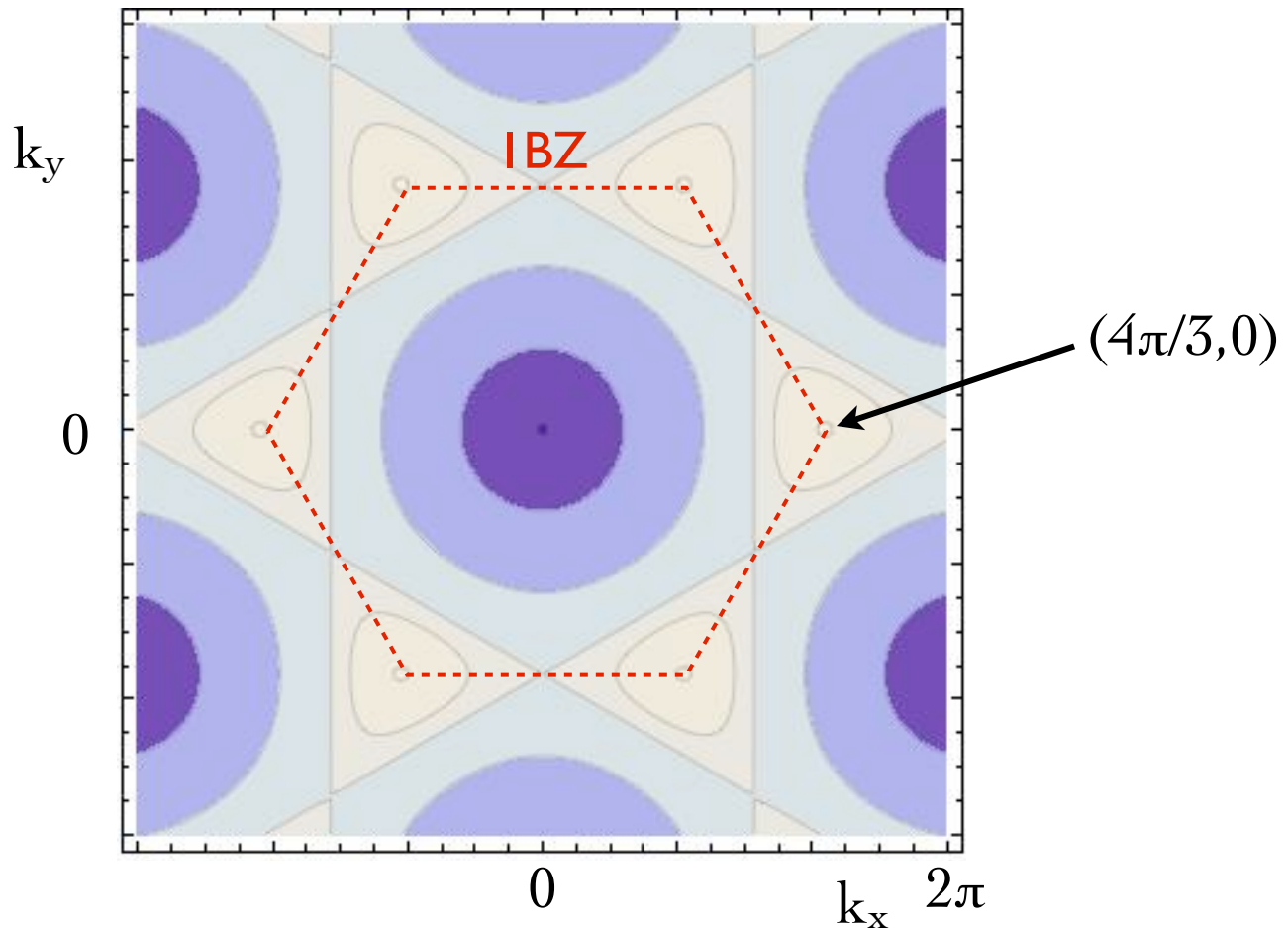
- Primitive vectors:  $(1, 0)$   $(-\frac{1}{2}, \frac{\sqrt{3}}{2})$   $(-\frac{1}{2}, -\frac{\sqrt{3}}{2})$

- Energy:

$$\epsilon(k) = \epsilon_0 - 2\gamma \left( \cos k_x + 2 \cos \frac{k_x}{2} \cos \frac{\sqrt{3}k_y}{2} \right)$$

# 2d hexagonal lattice

$$\epsilon(k) = \epsilon_0 - 2\gamma \left( \cos k_x + 2 \cos \frac{k_x}{2} \cos \frac{\sqrt{3}k_y}{2} \right)$$



# Graphene



- Single layers of honeycomb lattice of carbon
- First systematically exfoliated and studied by A. Geim + K. Novoselov, 2004.
- Nobel prize, 2011
- Interesting because it intrinsically has a *point Fermi surface*