

# Physics 23

## Assignment 3 Solutions

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### E25-4

a) The electro-static forces on the two masses are opposite and equal, so by Newton's Law we have

$$m_1 a_1 = m_2 a_2, \quad (1)$$

where  $a_1$  and  $a_2$  the magnitudes of  $m_1$ 's and  $m_2$ 's accelerations, respectively. Taking  $m_1 \approx 6.31 \times 10^7 kg$ , we find  $m_2 \approx 4.97 \times 10^{-7} kg$ .

b) Since the only force acting on the particles are electro-static, for force on one of the masses, for instance  $m_1$ , is

$$F_{on1} = m_1 a_1 = \frac{q^2}{4\pi\epsilon_0 r^2}, \quad (2)$$

where  $q$  is the common charge of the particles and  $r$  is the distance between them. Here all quantities are written as magnitudes of vectors since we have a simple 1-D problem. Putting in the numbers given to solve for  $q$ , we obtain  $q = 7.2 \times 10^{-11} C$ .

### E25-8

a) Define a Cartesian coordinate system with the  $+2q$  charge at the origin. The forces on  $+2q$  due to the three other charges are

$$\mathbf{F}_{+q \text{ on } +2q} = \frac{(2q)(q)(-\hat{\mathbf{j}})}{4\pi\epsilon_0 r_{+q,+2q}^2} = \frac{-\hat{\mathbf{j}}2q^2}{4\pi\epsilon_0 a^2} \quad (3)$$

$$\mathbf{F}_{-2q \text{ on } +2q} = \frac{(2q)(-2q)(-\hat{\mathbf{i}})}{4\pi\epsilon_0 r_{-2q,+2q}^2} = \frac{\hat{\mathbf{i}}4q^2}{4\pi\epsilon_0 a^2} \quad (4)$$

and

$$\mathbf{F}_{-q \text{ on } +2q} = \frac{(2q)(-q)(-\hat{\mathbf{i}}\cos(\pi/4) - \hat{\mathbf{j}}\sin(\pi/4))}{4\pi\epsilon_0 r_{-q,+2q}^2} = \frac{2^{1/2}(\hat{\mathbf{i}} + \hat{\mathbf{j}})q^2}{4\pi\epsilon_0 2a^2}. \quad (5)$$

Here the bold face denotes vectors while  $\hat{i}$  and  $\hat{j}$  are unit vectors in the x and y directions, respectively. The sine and cosine factors come from resolving the force due to -q into x and y components. The angle between the force vector in question and the x-axis is  $\pi/4$  radian.

We can then sum all the forces in the x-direction to obtain the horizontal component of the total force on +2q. We have

$$F_{on +2q,x} = \frac{4q^2}{4\pi\epsilon_0 a^2} + \frac{2^{1/2}q^2}{4\pi\epsilon_0 2a^2} \approx 2.34N. \quad (6)$$

Based on our coordinate system, this force is in the positive x-direction.

b) The vertical force on +2q has contributions from the y-components of the forces from +q and -q,

$$F_{on +2q,y} = -\frac{2q^2}{4\pi\epsilon_0 a^2} + \frac{2^{1/2}q^2}{4\pi\epsilon_0 2a^2} \approx -0.642N. \quad (7)$$

Note that +q charge applies a force on the +2q charge only in the y-direction, and the sign of its contribution is negative given our coordinate system.

### E25-16

First, define a coordinate system in which the rod lies completely on the x-axis with its ends at  $x=0$  and  $x=L$ . Let the distance between the right end of the rod and the charge +q be  $d$ . (the book used  $x$  which is confusing) We divide the rod into infinitely many differential (or very small) segments of length  $dx$  each and calculate the force on +q due to each bit of the rod. Finally, we obtain the total force on +q by summing up the contributions of all infinitesimal segments via integration across the length of the rod. The force on +q due to the differential amount of charge contained in length  $dx$  on the rod is

$$dF_{on+q} = \frac{qdQ}{4\pi\epsilon_0(L+d-x)^2} = \frac{dxQ/L}{4\pi\epsilon_0(L+d-x)^2}. \quad (8)$$

Here  $dQ = (Q/L)dx$  is the infinitesimal amount of charge carried on the small length  $dx$ . Note that the linear charge density of the rod is simply  $Q/L$  since the charge is uniformly distributed. With out coordinate system,  $L+d-x$  gives the distance between the differential bit of charge  $dQ$  and the observing charge, +q. We did not include vector symbols in the above expression since the problem is 1-D and it is clear that the total force on +q is in the positive x-direction.

Finally, we integrate Eqn.(8) across the length of the rod from 0 to  $L$  to obtain

$$F_{on+q} = \int dF = \int_0^L \frac{dxQ/L}{4\pi\epsilon_0(L+d-x)^2} = \frac{qQ}{4\pi\epsilon_0 L} \left( \frac{1}{d} - \frac{1}{d+L} \right). \quad (9)$$

**E25-22**

a) The  ${}^4\text{He}$  particle carries charge  $+2e \approx 2 \times 1.6 \times 10^{-19}\text{C}$  while the  ${}^{234}\text{Th}$  nuclei has charge  $+90e \approx 90 \times 1.6 \times 10^{-19}\text{C}$ . (since it has 90 protons) The magnitude of the repulsive force between them is given by Coloumb's Law,

$$F_{He,Th} = \frac{Q_{He}Q_{Th}}{4\pi\epsilon_0 R_{He,Th}} \approx 290\text{N}. \quad (10)$$

b) The only force on the He particle is the electro-static force, we thus have

$$F_{He,Th} = m_{He}a_{He}, \quad (11)$$

where  $m_{He} = 4m_{proton} \approx 4 \times 1.66 \times 10^{-27}\text{kg}$ . We find  $a_{He} \approx 4.4 \times 10^{28}\text{m/s}^2$ .

**P25-4**

a) In equilibrium, the net force on each ball is 0. In particular, the horizontal component of the tension pulling on each ball must cancel with the electro-static force (which is purely horizontal and opposite to the horizontal component of tension for a given ball). The electro-static force on each ball has magnitude

$$F_E = \frac{q^2}{4\pi\epsilon_0 x^2}. \quad (12)$$

The vertical component of the tension on each ball must balance gravity, we therefore have

$$T_y = T\cos(\theta) = mg. \quad (13)$$

Solving for the magnitude of the tension T and then computing its x-component, we obtain

$$T_x = T\sin(\theta) = mg\tan(\theta) \approx mg\sin(\theta) = \frac{mgx}{2L}. \quad (14)$$

Here we used the fact that  $\tan(\theta) \approx \sin(\theta)$  for small  $\theta$ . Equating  $T_x$  to  $F_E$  and solving for x yields the desired result.

b) Once again equate  $T_x$  to  $F_E$  but now solve for q using a given x. We obtain  $q \approx 2.28 \times 10^{-8}\text{C}$ .

**P25-8**

Initially, the charge -q experiences no net force since it is exactly half way between two charges of +Q at distance d apart. Now we move -q upwards a small distance y perpendicular

to the line joining the  $+Q$  charges (the x-axis, or  $y=0$ ). There will then be a net electrostatic force on  $-q$  in the negative y-direction (vertical), as the horizontal components of the forces due to the two  $+Q$  charges still cancel exactly. We now have a restoring net force that pulls  $-q$  back towards the line connecting the  $+Q$  charges. However, when  $-q$  reaches its initial position, it will have some finite velocity from being accelerated for some time by the restoring net force. We then expect  $-q$  to 'over shoot' and begin to experience a net force in the opposite ( $+y$ ) direction. The new net force will eventually stop the particle and turn it around back towards the positive y-direction. In effect,  $-q$  will undergo harmonic oscillation about the line joining the  $+Q$  charges. When  $-q$  is at a vertical distance  $y$  from the x-axis, the total electro static force on it is

$$F_{net} = -\frac{2qQ\sin(\phi)}{4\pi\epsilon_0 R_{-q,+Q}} = -\frac{2qQ\sin(\phi)}{4\pi\epsilon_0 [(d/2)^2 + y^2]} = -\frac{2qQy}{4\pi\epsilon_0 [(d/2)^2 + y^2]^{3/2}} \quad (15)$$

in the y direction. (+ or - y don't really matter here) Here  $\phi$  is the angle that position vector of  $-q$  makes with the x-axis when displaced vertically by  $y$  and we used

$$R_{-q,+Q} = [(d/2)^2 + y^2]^{1/2}, \quad (16)$$

so that

$$\sin(\phi) = \frac{y}{R_{-q,+Q}} = \frac{y}{[(d/2)^2 + y^2]^{1/2}}. \quad (17)$$

Next, note that  $y$  is assumed to be a very small number compared to  $d$  and so  $y^2$  is vanishingly small compared to  $(d/2)^2$ . We can then approximate  $F_{net}$  by dropping the  $y^2$  in the denominator. We then have

$$F_{net} \approx -\frac{2qQy}{4\pi\epsilon_0 (d/2)^3}. \quad (18)$$

Rewriting Eqn.(18) in a more enlightening form, we find

$$\frac{d^2y}{dt^2} \approx -\frac{4qQ}{m\pi\epsilon_0 d^3}y, \quad (19)$$

which is the differential equation for a simple harmonic oscillator with angular frequency

$$\omega^2 = \frac{4qQ}{m\pi\epsilon_0 d^3}. \quad (20)$$

Finally, taking the period  $T = 2\pi/\omega$  gives the desired result.

**P25-10**

a) Since all the  $Cs^+$  ions are arranged symmetrically at the same distance around the  $Cl^-$ , the net force on the Chlorine ion is zero.

b) After removing one  $Cs^+$ , note that there are still four  $Cs^+$  ions in symmetric positions around the Chlorine and so the forces due to those still cancel. Thus the net force on the  $Cl^-$  comes only from the remaining Cesium ion at  $d = (0.2^2 + 0.2^2 + 0.2^2)^{1/2}nm$  away. We then find

$$|\mathbf{F}_{net\ on\ Cl^-}| = \frac{e^2}{4\pi\epsilon d^2} \approx 1.92 \times 10^{-9}N. \quad (21)$$

Here  $e$  is the electron charge and the double vertical lines denote the magnitude of the net force vector.

**P25-11**

Let the line joining the two  $+q$  charges be on the x-axis. Without loss of generality, we can assume that the test charge  $Q$  is at a distance  $y=R$  above the x-axis. The net force on  $Q$  is in the y-direction ( $+$  or  $-$  does not matter here) since the horizontal components of the forces exerted by the  $+q$  charges cancel. The total force on  $Q$  is therefore

$$|\mathbf{F}_{on\ Q}| = \frac{2qQ\sin(\phi)}{4\pi\epsilon_0 d}, \quad (22)$$

where  $d$  is the distance between  $Q$  and either  $q$  and  $\phi$  is the angle between the position vector of  $Q$  and the x-axis. Next we use

$$d = (a^2 + R^2)^{1/2} \quad (23)$$

and

$$\sin(\phi) = \frac{R}{d} = \frac{R}{(a^2 + R^2)^{1/2}} \quad (24)$$

to obtain

$$|\mathbf{F}_{on\ Q}| = \frac{2qQR}{4\pi\epsilon_0(a^2 + R^2)^{3/2}}. \quad (25)$$

The value of  $R$  that maximizes  $|\mathbf{F}_{on\ Q}|$  is given by

$$\frac{d|\mathbf{F}_{on\ Q}|}{dR} = 0 = \frac{2qQ}{4\pi\epsilon_0} \left[ \frac{1}{(a^2 + R^2)^{3/2}} - \frac{3R^2}{(a^2 + R^2)^{5/2}} \right]. \quad (26)$$

The result  $R = \pm a/2^{1/2}$  immediately follows.